

# Supplementary Information for “Finding Density Functionals with Machine Learning”

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The purpose of this supplementary material is to provide the information necessary to reproduce our machine learning approximation (MLA) to the kinetic energy functional exactly, as well as extra details about the nature of the MLA. Table III gives the first 100 potential parameters  $a_i$ ,  $b_i$ , and  $c_i$  to double precision. Along with the equations in the main paper and details of Numerov’s method, one should be able to reproduce the corresponding weights  $a_j$  of the functional on line 4 of Table I in the main text.

## SELF-CONSISTENT DENSITIES

Table I gives the percentage of variance lost,  $(1 - \sum_{j=1}^{\ell} \lambda_j / \sum_{j=1}^G \lambda_j) \times 100\%$ , in taking the first  $\ell$  eigenvalues in the principle component analysis (PCA) projection. For each projection, there are  $G = 500$  eigenvalues, but only a few are necessary to represent the local variation in the densities. For example, with  $N = 1$ ,  $m = 30$  and  $\ell = 5$ , we retain 99.98% of the variance when projecting onto the PCA subspace. In Table II, we coarsely optimize the parameters  $m$  and  $\ell$  in the PCA projection.

$N\ell$	1	2	3	4	5	6	7	8
1	35	3	0.8	0.07	0.02	0.004	0.002	0.0003
2	45	15	3.7	0.36	0.10	0.019	0.006	0.0015
3	44	18	4.9	0.78	0.16	0.043	0.012	0.0037
4	48	24	9.5	1.9	0.40	0.087	0.023	0.0063

TABLE I. Percent of variance lost in taking the first  $\ell$  PCA eigenvalues with  $m = 30$ . Averaged over 100 PCA projections  $P_{m,\ell}(\mathbf{n})$  with randomly chosen centers  $\mathbf{n}$  in the test set.

## NUMEROV’S METHOD

We use Numerov’s method [1] to solve Schrödinger’s equation for  $N$  non-interacting spinless fermions confined to 1d box:

$$\left(-\frac{1}{2}\frac{\partial^2}{\partial x^2} + v(x)\right)\psi(x) = \epsilon\psi(x), \quad (1)$$

with boundary conditions  $\psi(0) = \psi(1) = 0$ . We discretize  $\psi(x)$  and  $v(x)$  on a uniform grid with spacing

$m\ell$	3	4	5	6
15	6.5	1.9	0.92	1.2
20	6.3	1.9	0.87	0.89
25	5.2	1.5	0.87	0.92
30	4.2	1.3	0.86	0.97
35	4.6	1.2	0.90	1.0
40	3.9	1.2	0.95	1.0
45	3.9	1.3	0.98	0.98

TABLE II. Mean absolute errors  $|\overline{\Delta T}|$ , in kcal/mol, on 100 randomly chosen self-consistent densities in the test set, averaged over  $N = 1, 2, 3$  and 4. This coarse optimization gives  $m = 30$ ,  $\ell = 5$ . Errors are less sensitive to  $m$  than  $\ell$ . For  $\ell \geq 7$ , the gradient descent search fails to converge in some cases.

$\Delta x = 1/(G - 1)$ :  $\psi_j = \psi(j/(G - 1))$ ,  $v_j = v(j/(G - 1))$  for  $j = 0, \dots, G - 1$ . Starting from  $\psi_0 = 0$  and  $\psi_1 = 1$ , we calculate the remaining  $\psi_j$  iteratively:

$$\psi_{j+1} = \frac{(2 - 5\Delta x^2 f_j/6)\psi_j - (1 + \Delta x^2 f_{j-1})/12\psi_{j-1}}{1 + \Delta x^2 f_{j+1}/12}, \quad (2)$$

where  $f_j = 2(\epsilon - v_j)$ . To determine the eigenvalues  $\epsilon^{(k)}$  and eigenfunctions  $\psi^{(k)}$ , we find the first  $N$  intervals that contain a root of  $\psi_{G-1}(\epsilon)$ , scanning from  $\epsilon = -3 \sum_{i=1}^3 a_i$  in steps of  $\Delta\epsilon = 1$ . For each interval  $k$ , we perform a binary search for the root, reducing the length of the interval to less than  $10^{-14}$ . The eigenvalue is taken as the midpoint of the interval (the error in  $\epsilon_k$  is less than  $5 \times 10^{-15}$ ). After normalizing the eigenfunctions  $\psi^{(k)}$ , the density and kinetic energy are given by:

$$n(x_j) = \sum_{k=1}^N \psi^{(k)}(x_j)^2, \quad T = \sum_{k=1}^N \epsilon_k - \sum_{j=1}^{G-1} n(x_j)v_j\Delta x. \quad (3)$$

[1] See e.g. E. Hairer, P. Nørsett, P. Syvert Paul and G. Wanner, *Solving ordinary differential equations I: Nons-tiff problems* (Springer, New York, 1993).

