



General terms	
DFT	density functional theory
G3	A standard database of molecular energy differences
GGA	generalized gradient approximation
GKS	generalized Kohn-Sham
HF	Hartree-Fock
KS	Kohn-Sham
LC20	A database of 20 lattice constants of solids
metaGGA	An extension of GGA that uses the KS kinetic energy
RPA	random-phase approximation
S22	A database of weakly bonded molecules
TDDFT	time-dependent density functional theory
XC	exchange-correlation
Approximate functionals	
B3LYP	The most popular functional in chemistry
B88	An exchange GGA that is the most popular in chemistry
LDA	local density approximation
PBE	A popular GGA in materials science
PW86	A predecessor of PBE
PW91	The GGA that PBE is a simplification of
SCAN	A recent meta-GGA

**Table 1** Table of acronyms in alphabetical order. References are given where they first appear in main text.

Almost all solve the Kohn-Sham (KS) equations (Kohn and Sham 1965) self-consistently, and use some approximation for the exchange-correlation (XC) energy as a functional of the (spin)-densities,  $E_{xc}[n]$ . This article surveys some of the more popular approximations. Some background is covered in Burke (2012), which is more focused on chemistry.

### 1.1 Commonly used approximations

The original approximation was suggested by Kohn and Sham themselves (Kohn and Sham 1965), namely the local density approximation (LDA), in which the XC energy density at each point in the system is replaced by that of a uniform electron gas with the density at that point. The exchange contribution was first written by Bloch (Bloch 1929), with correlation now well-known from quantum Monte Carlo simulations of the uniform gas (Ceperley and Alder 1980), parametrized in simple formulas (S. H. Vosko and Nusair 1980; Perdew and Wang 1992). LDA is remarkably accurate for geometries, but typically overbinds molecules by about 1 eV per bond, making it relatively useless for thermochemistry (Jones and Gunnarsson 1989).

The next step in complexity is the generalized gradient approximation (GGA), which creates an energy density using both the density and its gradient at each point (Burke et al 1997a). The basic concept, and its first real-

ization, were given already in the pioneering work of Ma and Brueckner (Ma and Brueckner 1968). This was carefully refined in the work of Langreth, leading to the Langreth-Mehl functional (Langreth and Mehl 1981, 1983). Perdew pioneered the use of real-space cutoffs to create GGA's, leading to the PW86 functional (Perdew and Wang (1986)). The highpoint of this detailed construction was the PW91 functional (Perdew 1991; J. P. Perdew and Filolhais 1992; Burke et al 1997b), and in 1993, it was shown that they yield useful accuracy for binding energies, i.e., errors of about 6-10 kcal/mol (1 eV = 23 kcal/mol) (Gill et al 1992). The most commonly used GGA in materials today is a simplification of the PW91 form called PBE (Perdew et al 1996a), while in chemistry it is BLYP, with B88 exchange (Becke 1988) and Lee-Yang-Parr correlation (Lee et al 1988). Both these are trained on Noble gas atomic energies, yielding more accurate energies for those atoms (Elliott and Burke (2009); Burke et al (2016)).

The last standard step is to create a hybrid of GGA with the exact exchange energy from a Hartree-Fock calculation, by replacing a fraction,  $a$ , of the GGA exchange with the Hartree-Fock (HF) exchange, as first suggested by Becke (Becke 1993). This fraction is 20% in the famous B3LYP functional, (which stands for Becke, 3-parameter, Lee-Yang-Parr), the most commonly used approximation in chemistry today (Becke 1993; Lee et al 1988; S. H. Vosko and Nusair 1980; P. J. Stephens and Frisch 1994). Its analog is PBE0 (Perdew et al 1996b; Ernzerhof and Scuseria 1999) which uses 25% mixing.

At least 80% of all DFT calculations currently being performed use one of the approximations mentioned above. One can think of more sophisticated approximations using more ingredients of the density, as in Jacob's ladder (Perdew and Schmidt (2001)). An important principle of progress in making density functional approximations is to ensure that (almost) everywhere, each level of approximation performs at least as well as the previous level, and also improves some features, making the previous level obsolete (except to save computational time). One should avoid having different functionals for different purposes.

## ***1.2 Beyond ground-state DFT***

What has been described above is generic ground-state DFT in the non-relativistic limit and for collinear magnetic fields. The single largest use of DFT beyond this domain is using time-dependent DFT (TDDFT) (Runge and Gross 1984; Maitra et al 2004; Maitra 2016) in the linear response regime to extract electronic transition frequencies (Casida 1996; M. Petersilka and Gross 1996). As much as 10% of all DFT publications include TDDFT estimates of excitations. Such calculations almost all use the adiabatic approximation for the XC kernel, and so are uniquely determined by

the choice of ground-state approximation. Most such calculations are for molecules (Burke et al 2005), as a non-local contribution is needed to yield details of bulk semiconductor spectra correctly (such as exciton binding energies), and that contribution is missing from the commonly-used functionals mentioned above (Onida et al 2002; Martin et al 2016).

Other extensions include magnetic DFT for non-collinear fields, relativistic DFT (Engel and Dreizler 2011) and even QED (Flick et al 2015), DFT at finite temperatures using the Mermin theorem (Smith et al 2017), coupling between nuclei and electrons (Requist and Gross 2016), and so on. Again, standard ground-state approximations are used unless a particular effect requires going beyond these, such as extracting double excitations (Maitra et al 2004).

## 2 Recent developments

This section is devoted to developments over the past twenty years.

### 2.1 RPA-type functionals

The famous adiabatic-connection fluctuation-dissipation formula of DFT yields  $E_{xc}$  in terms of the a frequency and coupling-constant and spatial integral over the density-density response function (Langreth and Perdew 1975; Harris and Jones 1974; Gunnarsson and Lundqvist 1976). This can be extracted directly from the KS response function, constructed from the occupied and unoccupied orbitals, and the XC kernel of TDDFT (Gross and Kohn 1985). Ignoring the kernel yields RPA, also known as direct random-phase approximation (RPA) (Furche 2001; Langreth and Perdew 1977, 1975; Chen et al 2017), also known as TD-Hartree. This scheme thus yields a fifth-rung (in Jacob's ladder) approximation that can be costly to evaluate, but the relative burden is always decreasing (Furche 2008; Eshuis et al 2010).

Direct RPA overcorrelates systems, because it includes only 'bubble' diagrams in the many-body expansion of the energy, and misses other contributions at higher-order that reduce correlation. It also has difficulties with self-interaction, because it yields finite correlation energies even for only one electron. These two effects yield inaccuracies in the dissociation energies of molecules. Recent progress has included various approximations to the XC kernel to yield improved energetics and computational cost-savings. A very recent development, using the exact frequency-dependent exchange kernel, and a clever and physically motivated resummation of higher-orders, appears to overcome stability problems and even allows the binding curve of  $N_2$  to be accurately calculated (Erhard et al 2016; Burke 2016) (which is

difficult even in coupled-cluster theory, due to the multireference nature of the wavefunction at large separations).

## 2.2 *Weak interactions*

Tremendous progress has been made in the last two decades for including weak, van der Waals interactions (vdW) into DFT calculations. The standard functionals model covalent, ionic, and metallic bonds reasonably well, but fail for longer, weaker bonds. Because of their semilocal nature, they cannot yield weak-binding that falls off as  $R^6$ , where  $R$  is the separation between two atoms or molecules. Thus corrections must be added to the standard functionals to capture these effects. There are three main schemes for doing this.

In the domain of explicit non-local density functionals, there is the sequence of approximations originally developed by Langreth and Lundqvist and co-workers (Dion et al 2004). These approximations are derived non-empirically, remarkably starting from contributions to the uniform gas correlation energy. The original was from 2004, and an improved parameterization was given in vdW2 (Klimeš et al 2009). These explicit density functionals can be applied to all materials, from molecules to solids, but have been designed assuming the system has a gap (not a metal) (Berland et al 2015).

At the extreme opposite end of the scale, there are the parametric schemes of Grimme and co-workers (Grimme 2006), commonly referred to as DFT-D. These are not explicit density functionals, but rather additional energies that include estimates of the  $C_6$  (and higher-order) coefficients in the asymptotic expansion of the vdW energy between atoms, combined with a damping factor to keep the contributions finite as the separation reduces. Such schemes require empirical parameters for each atom for a given standard XC approximation, but can yield highly accurate energy curves for small weakly bound molecules (Burns et al 2017).

In between is the scheme developed by Tkatchenko and Scheffler, which requires only one parameter for a given XC functional, and produces accurate add-on corrections to DFT energies (Tkatchenko and Scheffler 2009). This has been expanded to incorporate collective electrostatic effects so that metals and materials in many dimensions and on different length scales can be treated (Hermann et al 2017). A final scheme is that begun by Becke and Johnson, which uses the dipole moment of the exchange hole to determine  $C_6$  (and higher) coefficients (Johnson and Becke 2006).

Of course, more expensive treatments, such as RPA mentioned above, automatically include approximations to the VdW forces.

### 2.3 *meta-GGA's*

The third rung of Jacob's ladder is the meta-GGA, which adds a new ingredient beyond that of the density and its gradient. This is most often chosen to be the (positive) kinetic energy density of the KS orbitals. The aim for a good meta-GGA is to aim for the accuracy of hybrids without the computational overhead of the exact exchange contribution. The cost of exact exchange is relatively manageable in molecular calculations with atom-centered basis functions, but can often be prohibitive when using periodic boundary conditions and plane-wave basis sets. Running a range-separated hybrid (see next section) can sometimes take 1000 times longer than a typical GGA in a materials calculation.

Perdew and co-workers and many others have spent several decades developing meta-GGA's, with many flawed attempts (Sun et al 2015). But the most recent effort, called strongly constrained and appropriately normed semilocal density functional (SCAN), has passed many standard tests and appears very promising to join the pantheon of commonly-used functionals. The G3 dataset (L.A. Curtiss and Raghavachari 2005) is a standard set of chemical bonds that LDA overbinds by about 3 eV (typically about 1 eV/bond). PBE reduces this to about 1 eV, while SCAN reduces this to about 1/4 eV. SCAN also has errors that are 2-3 times smaller than PBE on the S22 dataset (Jurecka et al 2006) of weakly bonded systems. At the same time, SCAN reduces errors in lattice parameters on the LC20 data (Sun et al 2011) set from about 0.05Å in PBE to about 0.01Å. SCAN also yields better water properties than PBE (Perdew et al 2017). On the other hand, the underestimation of chemical barrier heights by PBE is only mildly improved, by about a factor of 30% (whereas hybrids are often 2-3 times better). Thus, for many properties, SCAN yields accuracies similar to hybrid functionals, but at a fraction of the computational cost (for materials codes).

### 2.4 *Range-separated hybrids*

The theory behind range separation is an exact one, developed first by Andreas Savin (Toulouse et al 2009; Savin 1996; Leininger et al 1997). One simply writes the Coulomb repulsion as a sum of a short-ranged contribution (decaying more rapidly than the inverse of the separation) and a long-ranged contribution, which has no Coulomb singularity at zero separation. One can then include one contribution as an interaction in some generalization of the KS equations, and have the redefined XC contribution accounting for the other. This is all formally exact, and exact XC functionals exist for such schemes (though they differ from their regular KS counterparts).

But a plethora of choices now await. The first is the length scale on which the range separation is performed, often denoted  $1/\omega$ . This is a continuous

parameter, and since approximations will be made to the corresponding  $E_{XC}$ , the results are sensitive to it. Just like the fraction of exact exchange in global hybrids, there is always a temptation to adjust it.

Furthermore, the separation into long- and short-range contributions can occur for the exchange contribution (most common), the correlation, or both. Finally, one may wish to treat the short-range contribution with an approximate functional or, in other situations, the long-range contribution! For example, a vexing problem in TDDFT is to accurately calculate charge-transfer excitations of well-separated donor-acceptor complexes. Including long-range Fock exchange exactly works very well for this problem (Stein et al (2009)). On the other hand, the very successful HSE06 functional (Heyd et al 2003) is a range-separated hybrid, in which the long-range exchange is treated with an approximate functional, but the short-range exchange is treated exactly, to model screening in an extended insulator (Janesko et al 2009). This particular hybrid typically yields accurate gaps for moderate-gap semiconductors and insulators (Perdew et al 2017).

## 2.5 Gaps of solids

A crucial failure of the standard approximations is their inability to predict gaps of semiconductors and insulators. The LDA underestimates the gap of bulk Si by a factor of 2, and makes Ge a metal, and GGA's fare little better. From very early on, a great strength of the GW method (Aryasetiawan and Gunnarsson 1998) has been its ability to provide accurate and reliable gaps.

To understand this issue, it is important to first note that the KS gap of a periodic solid does *not* match the fundamental gap of the solid (Perdew et al 1982; Perdew 1985). All indications suggest that in fact LDA and GGA yield reasonably accurate KS gaps (ie close to the KS gap found with the exact ground-state functional) (M. Grüning and Rubio 2006). But, unlike LDA and GGA, the exact functional allows access to the fundamental gap, which is just  $I - A$ , the difference between the ionization energy and electron affinity of a system. Consider a very large but finite cluster of material. One can then add and subtract an electron to find  $I$  and  $A$  and deduce the exact gap. In fact, modern methods exist for doing this in a periodic calculation (M. Stadele and Vogl 1999). But in LDA or GGA, the added electron or hole delocalizes over the entire system and, because of their lack of a derivative discontinuity,  $I - A$  collapses to the HOMO-LUMO KS energy difference, i.e., the KS gap.

A great success of the past two decades has been the accurate calculation of moderate gaps using hybrid functionals such as HSE06 (Heyd et al 2003). This is achieved by going to a generalized KS scheme (Seidl et al 1996), in which the orbital dependent part of the functional is treated as in HF theory, not pure KS theory (which would require treating it with optimized effective potential (OEP) methods (Kümmel and Kronik 2008)). By having

an orbital-dependent potential, one can show that the generalized Kohn-Sham (GKS) gap of such a calculation does match  $I - A$  and because of the orbital-dependence, an approximate derivative discontinuity is included. Thus the GKS gap in such a calculation is the approximate fundamental gap, and is wider than the corresponding KS gap. This is how hybrid functionals and meta-GGA's yield wider and generally more accurate gaps than GGA's (Perdew et al 2017).

### 3 Challenges and hopes

Here we review some of the more depressing failures of our current approximations.

#### 3.1 Strong correlation

The failure of DFT with standard approximations for strongly correlated systems has been known since before its invention, as HF has problems for  $H_2$  when it is stretched (Heitler and London 1927). The problem can be analyzed and related to localization/delocalization errors of the standard approximations when integer (or half-integer) numbers of electrons localize on different sites (Cohen et al 2008). This is often called static correlation in quantum chemistry and involves the KS gap between two states becoming very small, and the exact many-body wavefunction becoming an (almost) equal mixture of two Slater determinants.

It is important to note that the difficulties lie only with the failure of approximations under these circumstances, rather than the KS scheme itself. This can be beautifully illustrated with the two-site Hubbard model, for which it is trivial to construct the exact KS system, even when strongly correlated (Carrascal and Ferrer 2012). Even in realistic cases (albeit in 1d), the KS equations for a strongly correlated system always can be made to converge to the exact ground-energy and density if the exact XC functional is used (Wagner et al 2013). But of course many of the features of the KS system do not resemble those of the physical system under such conditions (Carrascal et al 2015). This point is often confused by practitioners of many-body theory. The differences between KS response functions and the many-body analogs for strongly-correlated systems is not a signal that a density functional approximation is failing to yield accurate energies for such systems.

But strong correlation in solids is even more difficult than static correlation in molecules. To see this simply, consider chains of uniformly spaced H atoms. As the spacing is increased, an electron localizes on each site. For  $H_2$ , the true wavefunction combines two Slater determinants. But for

$H_4$ , there are 4 such determinants, and a DFT calculation with, say LDA, will break symmetry into 4 different solutions, one of which will have lowest energy (Wagner et al 2014). As the number of atoms in the chain grows, so does the number of nearly degenerate solutions, separated by spin excitations of very small energy. In the thermodynamic limit, these become infinite, and the usual quantum chemical starting point of a single Slater determinant becomes hopeless (Qiu et al 2017). Because this is such an important problem, vast amounts of research have been performed studying this limit, especially by the group of Weitao Yang (Zheng et al 2011), but also by Scuseria (Motta et al 2017), and Becke (Johnson and Becke 2017).

### 3.2 *The role of empiricism*

The most practical systematic approach to the construction of density functionals has that been championed by Perdew: Combine exact conditions that are relevant to a given level of approximation with appropriate norms, such as the uniform gas or hydrogen atom, to create approximations of tremendous generality (Medvedev et al 2017). A key aspect of this approach is that, as one rises on Jacob's ladder, each successive approximation works better than the previous one (or at least is no worse) under almost all circumstances (and the cases where it does not are usually very informative). Thus, for a given computational cost, there is a single (or at least, very similar) obvious choice that is rarely worse than using a lower rung. Use of appropriate norms may appear empirical, but it can be understood as choosing parameters in the approximations to capture limits that have not been fully derived as yet, but in principle could be. The extremely successful B88 exchange functional (Becke 1988; Elliott and Burke 2009) can be viewed as incorporating appropriate norms (exchange energy of atoms).

On the other hand, the profusion of inexpensive computing resources has led to many databases with either experimental results or those of high-level computational chemistry against which new approximations can be tested. It has also led to empirical fitting of density functional approximations with many parameters, as championed by Truhlar and coworkers (Zhao and Truhlar 2008, 2006; Zhao et al 2006). This approach typically produces more accurate approximations than those of Perdew et al for the systems and properties fitted, and for related systems and properties. But it does not yield single universal approximations that generically improve over previous steps on the ladder. Such approximations can fail badly when applied beyond their range of applicability.

An entirely new approach to functional approximation is to use machine-learning to learn from accurate data (Snyder et al 2012; Rupp et al 2012; Bartók et al 2010). This differs from the earlier approaches, because it automatically includes highly nonlocal contributions, as captured for example

in the kernel which measures density separations by integrals over the entire system (Snyder et al 2012; Li et al 2016b; Vu et al 2015; Snyder et al 2013). Recent advances include the first KS-MD simulation with an ML-DFT approximation to bypass solving the KS equations (Brockherde et al 2017), and accurate approximation of the full interacting functional (including XC) (Li et al 2016a) even for strong correlation and even for extended systems (but so far, only in 1d, because of the cost of generating accurate data).

### 3.3 *New horizons*

An entirely new arena for DFT has grown immensely in the last decade or so is applications to warm dense matter (Smith et al 2017), with temperatures significant on the electronic scale (about 100,000K), but not so high that Thomas-Fermi theory (or classical behavior) dominates. Applications range from modeling planetary interiors to inertial confinement fusion. This field is so 'hot' that even the input to thermal LDA, the XC energy of a uniform gas as a function of temperature, is only now being calculated at high accuracy (Groth et al 2017).

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